

AROUND SMALE'S 14TH PROBLEM

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ABSTRACT. We consider the Lorenz equations on \mathbb{R}^3 . We prove analytically the existence of T-parameters, where there exists a heteroclinic connection between the singularities forming an invariant trefoil knot. Furthermore, we show that at a T-parameter the geometric model is a minimal representative for the equations, implying e.g. that the Lorenz equations contain all Lorenz knots (defined by the geometric model) as periodic orbits.

1. INTRODUCTION

We study the classical Lorenz equations as given in [Lor63]:

$$(1) \quad \begin{cases} \dot{x}(t) &= \sigma(y - x) \\ \dot{y}(t) &= \rho x - y - xz \\ \dot{z}(t) &= xy - \beta z \end{cases}$$

Comment You should define notation used either here or in a notation at the start. Include also assumptions. Then state applications where problems like this arise.

This system has been extensively studied over the last decades, being the first system to demonstrate chaos in a deterministic low dimensional setting. State what is meant by "low dimensional" and what range of dimensions you are considering. How many components can functions x, y , and z have for example? It has been studied especially at the classical parameter values $\rho = 28, \sigma = 10, \beta = \frac{8}{3}$ which were those originally studied by Lorenz.

At these parameter values the three-dimensional motion given by the Lorenz equations (1) converges almost always onto its famous butterfly shaped strange attractor. Define "strange attractor". Although this is easily seen to be the case by running any ODE solver there is no analytical proof for it, and in general the equation has proven to be hard to study both analytically and numerically.

For these reasons, a simpler model, called the geometric Lorenz model, was developed [GW79, ABS77]. This model shares much of the characteristic of the equations as it also possesses a butterfly shaped strange attractor. At the same time, the geometric model is susceptible to analytic study and it can be proven to be chaotic

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in a specific sense, for example its attractor has a central limit theorem [HM07] and exponential decay of correlations [AM16]. The periodic orbits in the geometric model, which are called “Lorenz knots” have been studied topologically by Williams [Wil79] and Birman and Williams [BW83] (see also see also [Deh11, BK09]) and have a number of surprising knot properties. Define or explain ”knot properties” briefly.

Smale’s 14th problem [Sma00] asks whether the original system at the classical parameters can be proven to be topologically equivalent to the geometric model. What does classical mean in this context? This has been answered by Tucker [Tuc99], via a rigorous numerical proof. Tucker’s proof implies in particular that the Lorenz equations indeed possess a butterfly strange attractor.

An astounding result of Ghys [Ghy07] is that the set of Lorenz knots also arise as the set of periodic orbits for the geodesic flow on the modular surface. The modular flow is also a chaotic flow but a very different one. It is a hyperbolic flow arising from arithmetics Is this a typo or jargon? If the latter, define/explain, for which statistical properties can more readily be proven. It is defined on the complement of the trefoil knot, which is the simplest nontrivial knot. It is well known by numerical studies that for the Lorenz equations there exist points in parameter space, called T-points, for which there exists a heteroclinic connection connecting its three singular points [PI80, AF85, Byk80]. It has been suggested by the author in [Pin17] that this heteroclinic connection can be continued to an invariant knot passing also through infinity. For the first T-point at $\rho_0 \approx 30.8680$, $\sigma_0 \approx 10.1673$ and $\beta_0 = \frac{8}{3}$, the invariant knot is a trefoil knot Define trefoil, depicted in Figure 1

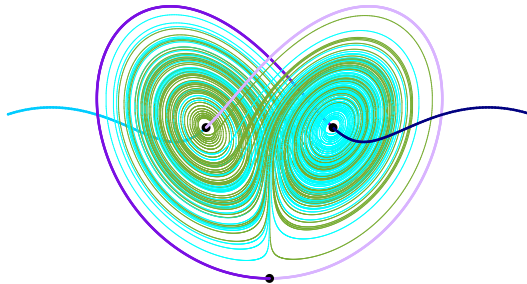


FIGURE 1. The trefoil knot at the first T-point, see also [CKO15].

In an attempt to explain the relation between the equations and the geometric model and the modular geodesic flow, Christian Bonatti and the author have constructed an extension of the geometric model [BP21], which is defined on the entire

three dimensional sphere, and for some parameter has an invariant trefoil shaped heteroclinic connection.

Our main theorem is an approximate answer to Smale's problem. It states that there exists a parameter for which the Lorenz equations can be continuously deformed to the geometric model. Furthermore the deformation, although not an equivalence, cannot create new periodic orbits.

Theorem 1.1. *There exists a parameter value for the Lorenz equations for which a trefoil shaped heteroclinic connection exists. At this parameter value, the extended geometric model is isotopic to the Lorenz equations.*

The main tool is the existence of a convenient cross section for the flow, which allows us to establish the return map corresponds to a symbolic dynamics. Define this term. on two symbols. Note that an isotopy is much weaker than a topological equivalence, as it can potentially lose much of the dynamics of the flow. Our second result is thus that this doesn't happen, namely, that the isotopy taking the geometric model to the Lorenz flow cannot kill any of its periodic orbits. This is done using the section in order to apply methods from surface dynamics [Thu88, Boy94, AF83].

Theorem 1.2. *Any periodic orbit in the geometric model is isotopic to a periodic orbit for the Lorenz equations at the T -point.*

An immediate corollary of our results above and Ghys' theorem [Ghy07] is thus the following.

Corollary 1.3. *The geodesic flow on the modular surface is the minimal representative to the Lorenz equation at the trefoil T -point. Namely, a periodic geodesic is isotopic to a periodic orbit of the Lorenz equations.*

Note that at the classical parameters it follows from Tucker's proof that any periodic orbit for the Lorenz equation is isotopic to a modular geodesic, but not every geodesic appears. The set that does appear is determined by the kneading sequence at that point. Define "kneading sequence".

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2. THE EXISTENCE OF A HETEROCLINIC TREFOIL

In this section we prove the existence of a trefoil heteroclinic connection passing through all four singular points for some point in the parameter space. The first step in the proof is establishing the existence of a global cross section, and the second is a winding number argument.

The cross section we use is modified from the one commonly used in numerical studies which has a constant z coordinate, and is reminiscent of Sparrow's use of the maximal values of the z coordinate of the orbit. It is a section also at parameter values for which the regular section fails to be one.

Proposition 2.1. *There exists an open simply connected domain A of parameters, for which for any point $\kappa = (\beta, \sigma, \rho)$ there exists a two dimensional topological rectangle $R_\kappa \subset \mathbb{R}^3$ with interior transverse to the κ -Lorenz flow, so that the forward orbit of any point in \mathbb{R}^3 that does not limit onto one of the singular points meets R_κ .*

Proof. We start by considering the hyperbolic paraboloid

$$\mathcal{P} = \{xy = \beta z\}$$

Which is the set of points for which $\dot{z} = 0$, containing the three singular points. The paraboloid \mathcal{P} divides \mathbb{R}^3 into two components, we call the part of $\mathbb{R}^3 \setminus \mathcal{P}$ above the origin the inside and the other component the outside.

For any orbit, its z coordinate decreases while it is on the inside, and increases if and only if it crosses \mathcal{P} to be on its outside.

A normal to \mathcal{P} is given by

$$N = (y, x, -\beta).$$

Taking its inner product with the vector field $X = (X_1, X_2, X_3)$, along the paraboloid \mathcal{P} we obtain

$$N \cdot X = \sigma y^2 - \sigma xy + \rho x^2 - xy - \frac{1}{\beta} x^3 y$$

The equation $N \cdot X = 0$ is quadratic in y , and thus the set of points $N \cdot X = 0$ for which the vector field is tangent to \mathcal{P} consists of a point at the origin, and of two smooth one dimensional curves which we denote by δ_- and δ_+ each containing one of the wing centers.

When $\rho > \frac{(\sigma+1)^2}{4\sigma}$ the two dimensional curves are far from the origin and the point at the origin is an isolated solution, while for smaller ρ the two curves have cusps touching the origin. Thus we will restrict our attention from now on to the domain

$$A = \left\{ \beta, \sigma, \rho \mid \beta, \sigma > 0, \rho > 1, \rho > \frac{(\sigma+1)^2}{4\sigma} \right\}.$$

Note that the domain A is simply connected (i.e. any loop in A can be retracted to a point). For any parameter in A , the orbits cross \mathcal{P} to its outside below the two curves δ_- and δ_+ and enter the inside above these curves. An example of the regions of entrance to the inside of \mathcal{P} and exit from it are depicted in Figure 2.

Recall that a large enough ellipsoid around the origin of the form

$$S = rx^2 + \sigma y^2 + \sigma(z - 2r)^2$$

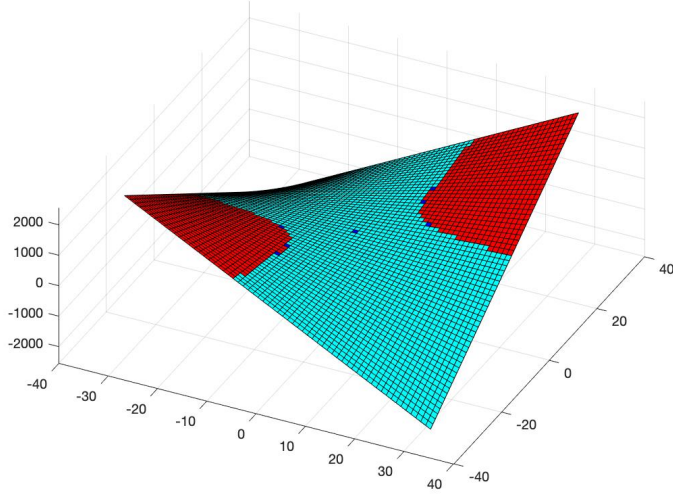


FIGURE 2. The paraboloid \mathcal{P} for the classical parameters, with the regions where the flow is directed upwards depicted in red and downwards flow is depicted in blue.

is transverse to the vector field X , so that orbits only enter the region bounded by the ellipsoid (see [Lor63, Spa82]). The curves δ_- and δ_+ intersect the ellipsoid in four points.

We next define another pair of arcs, α_1 and α_2 . First, fix an $\varepsilon > 0$. Define an arc continuing from each intersection point of $\delta_{\pm} \cap S$ downwards, in the intersection $\mathcal{P} \cap S$, until it reaches the plane $z = \varepsilon$. Then, connect the endpoints in the intersection of the plane and the ellipsoid (and in the inside of \mathcal{P}). The union of two of the downwards arcs and a horizontal arc is an arc α_i , $i = 1, 2$ so that the four arcs $\delta_-, \delta_+, \alpha_1, \alpha_2$ bound a closed rectangle R , contained in the union of \mathcal{P} and the plane $z = \varepsilon$.

The Lorenz equation have a symmetry of rotation about the z axis $\pi : (x, y, z) \mapsto (-x, -y, z)$. By construction, R is preserved by the symmetry and $\pi(\alpha_1) = \alpha_2$, $\pi(\delta_-) = \delta_+$.

The part of R that is contained in \mathcal{P} is disjoint from the set $N \cdot X = 0$, and therefore the flow is transverse to R on this set. The other part of R is contained in the intersection of the plane $z = \varepsilon$ with the inside of \mathcal{P} . Thus, the flow there satisfies $\dot{z} < 0$, and as this part of the rectangle is horizontal, the flow is transverse to R there as well.

Orbits cannot escape to infinity as close to infinity the flow is pointing inwards through a family of concentric large enough ellipsoids [Lor63]. It follows from the

shape of \mathcal{P} that there are no nonsingular orbits with constant z coordinate, the forward orbit of any point $x \in \mathbb{R}^3$ either has at some point an increasing and then decreasing z coordinate, in which case it must meet R , or the orbit limits onto a singular point (or is fixed at a singular point). This completes the proof. \square

Lemma 2.2. *The component of the intersection of the stable manifold and R containing the point on the z axis is a one dimensional curve dividing R into two components, yielding a symbolic coding on two symbols.*

Proof. Recall that the stable manifold of the origin is two dimensional and contains the z axis, thus it intersects R at least at one point, i.e. the point at which the z axis intersects the plane $z = \varepsilon$. From the transversality of the flow to R it follows, since the flow is tangent to the stable manifold, that the intersection of the stable manifold and R is transversal. The stable manifold is topologically a disk. It cannot have a periodic orbit on its boundary as in this case it would follow from the symmetry that it has two boundary components, in contradiction. Thus, the intersection is a one-dimensional manifold that connects to infinity on both sides.

The intersection curve η is preserved by the symmetry π , and does not pass through the wing centers p^+ and p^- . Thus it divides R into two symmetric subsets, each containing one of the wing centers.

Denote the Poincaré return map by φ . The curve η yields a natural way to endow the system (R, φ) with a symbolic dynamics with two symbols, corresponding to the two components of $R \setminus \eta$. \square

Denote by A and B the two components of $R \setminus \eta$. Each component has one of the wing centers on its boundary, and thus the transition A to A appears in the dynamics, and also B to B . This suggests the return map has the shape of two triangles. However, there is still too much freedom for this symbolic dynamics to be useful. This is why we next focus on special parameter values allowing us to pinpoint the behaviour of the return map.

2.1. Existence of the heteroclinic trefoil.

Proposition 2.3. *For each of the wing centers p^\pm , one side of its stable manifold connects it to infinity.*

This has been observed numerically, see for example [Spa82] and [CKO15].

Proof. Consider the quadrant

$$\mathcal{Q}_+ = \{(x, y, z) \mid x \geq x_0, y \leq x\},$$

where $x_0 = y_0 = \sqrt{\beta(\rho - 1)}$ is the coordinates of the fixed point p^+ . By considering the linearization at the fixed point one finds that there is always at least one negative

real eigenvalue λ , corresponding to an eigenvector v with $v_1 \neq 0$. Consider a vector v corresponding to λ with $v_1 > 0$. By considering the equation for \dot{y} it must have $v_2 \leq v_1$ and thus the vector v (based at p^+) points to the interior of \mathcal{Q} .

Next, consider $\partial\mathcal{Q}_+$. The flow is transverse to the half plane $x \equiv x_0$ when $y \leq x$, pointing outwards at each point. It points outwards through the plane $x = y$ whenever $z < \rho - 1$, and points inwards whenever $z > \rho - 1$.

In addition to $\partial\mathcal{Q}_+$ we may consider the ledge sticking into the interior of \mathcal{Q}_+ consisting of a wedge with one edge being the line $\{x = y, z = \rho - 1\}$ and a second edge on the half of δ_- within Q (note that these two curves meet at p^+), so that the interior of the ledge lies entirely outside \mathcal{P} as in the figure.

consider the half of δ_+ and its image under the flow. locally near p^+ each half of δ_+ flows, forming a two dimensional spiral with a cone point at p^+ . Its image cannot intersect the original curve δ_+ as then there would exist a compact region in space, composed of the spiral shell of δ_+ and its flowlines, and the union of the ledge and $\partial\mathcal{Q}_+$, from which the flow only exists, contradicting the fact the Lorenz flow has negative divergence everywhere.

Taking into account the flow directions on $\partial\mathcal{Q}_+$, one thus concludes that the image of the half of δ_+ contained there may either wind around δ_+ and intersect the ledge, or it may exit \mathcal{Q} through one of the faces where the flow points outwards. In both cases, the union of the shell of flowlines $\bigcup_{t \geq 0} \psi^t(\delta_-)$ together with the ledge and the outflowing part of $\partial\mathcal{Q}_+$ bounds an open cone with tip at p^+ , to which the flow is either tangent or exiting. Therefore, a component of the stable manifold of p^+ must be contained in this cone. As flowlines can enter this cone only from infinity (passing through the sphere S), this half of the stable manifold connects p^+ to infinity without intersecting the cross section R .

The same claim for p^- follows from the symmetry. □

Theorem 2.4. *There exists a point in the parameter space at which there exists a heteroclinic trefoil for the Lorenz equations.*

The main tool of this proof is the following theorem, proving existence of homoclinic orbits and their linking with the vertical lines

$$L^\pm = \{x = y = \pm\sqrt{\beta(\rho - 1)}\}$$

through the two wing centers.

Theorem 2.5 (Chen, Theorem 1.1 and Lemma 4.3 of [Che96]). *For any given positive number β and non-negative integer N , there exists a large positive constant $R_0(N, \beta)$ such that for each $R > R_0$ there is a positive number $s = s(N, \beta, R) = (2\beta + 1)/3 + O(R^{-1/2})$ such that the Lorenz system has a homoclinic orbit associated with the origin which rotates around L^+ exactly $[\frac{N+1}{2}]$ times, and rotates $[\frac{N-1}{2}]$ times around L^- .*

Proof. We define a family of isotopic loops l_κ^+ and l_κ^- each passing through infinity and through one of the wing centers, as follows. One half of the loop is the half of $\mathcal{W}^s(p^\pm)$ that connects p^\pm to infinity by Proposition 2.3. Now consider the image of half of δ_\pm that is on the other side of p^\pm (and is not in the quadrant considered in the proposition). For a small enough time segment this image is either on the outside of \mathcal{P} , including the region of \mathcal{P} where the flow flows outwards, or on the inside of \mathcal{P} (including the inward flowing region).

At parameters at which the image is inside \mathcal{P} , we choose l_κ^\pm to be the union of this half of δ^\pm and the half of $\mathcal{W}^s(p^\pm)$ that connects p^\pm to infinity by Proposition 2.3.

At parameters at which the image is outside \mathcal{P} , the half of δ^\pm and the flowlines emanating from it, until the first point they hit \mathcal{P} in its incoming region enclose together with \mathcal{P} a three dimensional region from which the flow can only exit, proving as in Proposition 2.3 that the other half of the stable manifold of p^\pm connects to infinity as well (and in particular, the two curves are isotopic). For these parameters we choose l_κ^\pm to be equal to the stable manifold $\mathcal{W}^s(p^\pm)$. In both of these cases, the choice of l_κ^\pm ensures that orbits arriving from inside \mathcal{P} can hit l_κ^\pm only at p^\pm .

The description in Theorem 2.5 of the homoclinic orbits implies that the positive half of the simplest such orbit links once with L^+ and does not link with L^- . It follows that the separatrix, i.e. the unstable manifold of the origin, at that parameter returns to \mathcal{P} from the $x, y > 0$ quadrant, linking δ^+ , and likewise l_κ^+ , once, and then connects to the origin after hitting R a single time. In particular it is not linked with l_κ^- . At the second homoclinic loop the separatrix links L^+ like in the first loop, and then continues to also link once with l_κ^- . It then connects to the origin after it hits R the second time.

Choose ε small enough so that the ball $B(\varepsilon)$ of radius ε about the origin does not intersect l_κ^\pm , and so that once the homoclinic loops enter the ball they do not leave it.

Consider the map ψ_T taking a point κ in parameter space to the point along the separatrix, (on its half γ^+ starting with positive x coordinate) to the point reached at time T when starting at distance $\varepsilon/2$ from the origin along γ^+ . ψ_T is continuous as the separatrix depends continuously on the parameters. This is true as it is a fixed point of a contracting operator on a Banach space depending continuously, in the C^1 topology, on the parameters of the ordinary differential equation.

Consider a path $\kappa(s)$ connecting two points in parameter space $\kappa(0)$ corresponding to the first homoclinic loop and $\kappa(1)$ to the second. Suppose for each point along $\kappa(s)$ there exists a finite time $T(s)$ for which $\psi_T(s)(\kappa(s)) \in B(\varepsilon)$.

The map $s \mapsto T(s)$ is continuous and thus the topological disk $D = \{(s, T) | 0 \leq s \leq 1, 0 \leq T \leq T(s)\}$ is mapped by $\psi_T(\kappa(s))$ to a disk in S^3 , with boundary contained in the union of the two different homoclinic loops and of $B(\varepsilon)$.

The disks' boundary ∂D winds once along l_κ^- and hence there exists a point (s_*, T_*) within the disk so that $\psi_{T_*}(\kappa(s_*))$ lies on l_κ^+ . The choice of l_κ^+ and l_κ^- ensures that any flowline hitting them not at p^\pm arrives from infinity without crossing R or \mathcal{P} . This implies that $\psi_{T_*}(\kappa(s_*)) = p^-$, and this is impossible for a finite time T_* . Therefore, the separatrix orbit for s_* cannot return to $B(\varepsilon)$ in a finite time for any s .

Therefore, for any such path there must exist a value of s for which the separatrix orbit does not return to $B(\varepsilon)$ for any finite time. Assume the separatrix orbit hits the cross section $R(\kappa(s))$ for any $s \in [0, 1]$ in a finite time $T(s)$. The path $\eta(s) = \psi_{T(s)}(\kappa(s))$ in R is continuous. In this case the ball $B(\varepsilon)$ can be stretched through the stable manifold of the origin and along $\eta(s)$, remaining topologically a ball. This forces every separatrix orbit to return to it along the way, in contradiction to the previous claim. Thus we conclude that there exists a parameter along $\eta(s)$ for which the separatrix does not hit R for any finite time. It follows that the separatrix there reaches p^- and is a heteroclinic orbit. Hence, there exists a parameter for which we have a heteroclinic connection along any path connecting the homoclinic orbits.

Let κ_* be such a parameter. It follows from the proof that the separatrix does not hit R before reaching p^- , and p^+ on its other side. Now it follows from 2.3 that the other half of each stable manifold connects directly to infinity and therefore the resulting knot is a trefoil knot as required. \square

Remark 2.6. *It follows that a heteroclinic trefoil exists at any fixed β , as the parameter domain A restricted to any fixed β is still simply connected.*

3. THE DYNAMICS AT A TREFOIL T-POINT

We next prove Theorem 1.2, showing that at a trefoil T-point, every Lorenz knot (and equivalently by Ghys any modular knot) appears as a periodic orbit for the Lorenz system.

Proof. Consider for a T-parameter where the heteroclinic connection is a trefoil knot the cross section R , divided into two components A and B by the stable manifold of the origin, and the return map φ as developed in previous sections.

The image of A under φ is a topological disk that includes both wing centers on its boundary: The wing center in A as it is a fixed point and thus is equal to its image, and the wing center in B as it is the limit of points that are adjacent to the stable manifold of the origin, and thus after approaching the origin they continue along the separatrix that connects to this fixed point.

Thus, the image of R under the return map φ has the shape of two bananas sharing the two fixed points p^\pm on their boundaries at their "corners". To define a symbolic dynamics, send a point x in R to the infinite sequence of letters A's and B's $(\dots, s(f^{-1}(x)), s(x), s(f(x)), \dots)$ where $s(x)$ is A if $x \in A$ and B if $x \in B$. As each

image crosses R from side to side, The map is onto, i.e. any symbol corresponds to at least one point in R .

The map need not be injective, however by Brouwer fixed point theorem, every periodic symbol corresponds to at least one point in R of a periodic orbit with the same period (under the first return map). This orbit is isotopic to a Lorenz knot as the images of A and B do not intersect and thus the orbit will always pass from A to B from above, and from B to A from below.

The extended geometric model in [BP21] corresponds to the full set of possible symbols in A and B , and any possible periodic symbol corresponds to a periodic orbit in the geometric model which is a Lorenz knot by the equivalence proven in [BP21] between the extended model and the classical Lorenz geometric model. \square

Remark 3.1. *Using the topological approach presented here to give an analytical solution to the problem will require a way to prove the nonexistence of attracting periodic orbits for the Lorenz equations, showing the above symbolic partition is generating. This seems to be hard as it is a local phenomenon that cannot be obstructed topologically. It is numerically observed that there are parameters for which such orbit appear, but at ρ values ~ 100 .*

REFERENCES

- [ABS77] V. S. Afraimovich, V. V. Bykov, and L. P. Shilnikov. On the origin and structure of the Lorenz attractor. *Akademiia Nauk SSSR Doklady*, 234:336–339, May 1977.
- [AF83] D. Asimov and J. Franks. Unremovable closed orbits. In *Geometric dynamics (Rio de Janeiro, 1981)*, volume 1007 of *Lecture Notes in Math.*, pages 22–29. Springer, Berlin, 1983.
- [AF85] K. H. Alfsen and Jan Frøyland. Systematics of the Lorenz model at $\sigma=10$. *Physica Scripta*, 31(1):15–20, 1985.
- [AM16] Vitor Araújo and Ian Melbourne. Exponential decay of correlations for nonuniformly hyperbolic flows with a $C^{1+\alpha}$ stable foliation, including the classical Lorenz attractor. *Ann. Henri Poincaré*, 17(11):2975–3004, 2016.
- [BK09] Joan Birman and Ilya Kofman. A new twist on Lorenz links. *J. Topol.*, 2(2):227–248, 2009.
- [Boy94] Philip Boyland. Topological methods in surface dynamics. *Topology Appl.*, 58(3):223–298, 1994.
- [BP21] Christian Bonatti and Tali Pinsky. Lorenz attractors and the modular surface. *Nonlinearity*, 34(6):4315–4331, 2021.
- [BW83] Joan S. Birman and R. F. Williams. Knotted periodic orbits in dynamical systems. I. Lorenz’s equations. *Topology*, 22(1):47–82, 1983.

- [Byk80] V. V. Bykov. Bifurcations of dynamical systems close to systems with a separatrix contour containing a saddle-focus. In *Methods of the qualitative theory of differential equations*, pages 44–72. Gor'kov. Gos. Univ., Gorki, 1980.
- [Che96] Xinfu Chen. Lorenz equations. II. “Randomly” rotated homoclinic orbits and chaotic trajectories. *Discrete Contin. Dynam. Systems*, 2(1):121–140, 1996.
- [CKO15] Jennifer L. Creaser, Bernd Krauskopf, and Hinke M. Osinga. α -flips and T-points in the Lorenz system. *Nonlinearity*, 28(3):R39, 2015.
- [Deh11] Pierre Dehornoy. Les nœuds de Lorenz. *Enseign. Math. (2)*, 57(3-4):211–280, 2011.
- [Ghy07] Étienne Ghys. Knots and dynamics. In *International Congress of Mathematicians. Vol. I*, pages 247–277. Eur. Math. Soc., Zürich, 2007.
- [GW79] John Guckenheimer and Robert F. Williams. Structural stability of lorenz attractors. *Publications Mathématiques de l’IHÉS*, 50:59–72, 1979.
- [HM07] M Holland and I Melbourne. Central limit theorems and invariance principles for lorenz attractors. *JOURNAL OF THE LONDON MATHEMATICAL SOCIETY-SECOND SERIES*, 76:345 – 364, 2007.
- [Lor63] Edward N Lorenz. Deterministic nonperiodic flow. *Journal of the atmospheric sciences*, 20(2):130–141, 1963.
- [PI80] N V Petrovskaya and Yudovich V I. Homoclinic loops of the saltzman-lorenz system. *Methods of Qualitative Theory of Differential Equations*, pages 73–83, 1980.
- [Pin17] Tali Pinsky. On the topology of the lorenz system. *Proceedings of the Royal Society A: Mathematical, Physical and Engineering Sciences*, 473(2205):20170374, 2017.
- [Sma00] Steve Smale. Mathematical problems for the next century. In *Mathematics: frontiers and perspectives*, pages 271–294. Amer. Math. Soc., Providence, RI, 2000.
- [Spa82] Colin Sparrow. *The Lorenz equations: bifurcations, chaos, and strange attractors*, volume 41 of *Applied Mathematical Sciences*. Springer-Verlag, New York-Berlin, 1982.
- [Thu88] William P. Thurston. On the geometry and dynamics of diffeomorphisms of surfaces. *Bull. Amer. Math. Soc. (N.S.)*, 19(2):417–431, 1988.
- [Tuc99] Warwick Tucker. The Lorenz attractor exists. *C. R. Acad. Sci. Paris Sér. I Math.*, 328(12):1197–1202, 1999.
- [Wil79] R. F. Williams. The structure of Lorenz attractors. *Inst. Hautes Études Sci. Publ. Math.*, (50):73–99, 1979.

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